



# A novel $\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$ microwave dielectric ceramic with ultra-low firing temperature

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## ABSTRACT

A new type of  $\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$  (CZM) ceramic was successfully prepared through the solid-phase reaction process at low sintering temperature. All the studied sintered samples exhibited a trigonal phase with space group of R-3c (No.167) by means of X-ray diffraction. Rietveld refinement analysis results further explored the crystalline structure of CZM. According to Scanning electron microscopy images, dense uniform surface morphology was observed at 525–625 °C for CZM ceramics. The correlations between structural characteristics and microwave dielectric properties of CZM ceramics were systematically investigated by calculating some chemical bond parameters on the basis of the chemical bonds theory of complex crystals. The analysis results of far-infrared reflectivity spectroscopy revealed that the absorption of structural phonon vibration at infrared frequencies played a dominant role in dielectric polarization contribution for CZM ceramics. Furthermore, the CZM sample sintered at 575 °C exhibited optimum microwave dielectric properties of  $\epsilon_r = 10.69$ ,  $Qf = 19,062$  GHz and  $\tau_f = -1.29$  ppm/°C.

## 1. Introduction

Because of the importance of integration and miniaturization for microwave circuit system, the low-temperature co-fired ceramics (LTCC) technology has been explored. Ceramics with high  $\epsilon_r$  values, high  $Qf$  values, near-zero  $\tau_f$  values and low sintering temperatures are widely tried to design the LTCC devices [1,2]. It is urgently to develop microwave dielectric ceramics with low sintering temperatures and excellent microwave dielectric properties to satisfy the rapid advancement in communication industry. Therefore, there is a strong drive to develop new microwave dielectric ceramics [3–8], lower temperatures by adding sintering additives [9–11] and improve microwave dielectric properties by ion substitution [12–15]. Especially, a series of new microwave dielectric ceramics, such as  $\text{MoO}_3$ ,  $[\text{Ca}_{0.55}(\text{Sm}_{1-x}\text{Bi}_x)_{0.3}]\text{MoO}_4$  and  $(\text{Ca}, \text{Bi})(\text{Mo}, \text{V})\text{O}_4$  ceramics, have been studied in detail [16–18].

Recently, as shown in Table 1, Mo-based microwave dielectric ceramics, such as  $\text{Sm}_2\text{Zr}_3(\text{MoO}_4)_9$ ,  $\text{Nd}_2\text{Zr}_3(\text{MoO}_4)_9$  and  $\text{Eu}_2\text{Zr}_3(\text{MoO}_4)_9$ , etc., have achieved great attention owing to the excellent microwave dielectric properties and low sintering temperatures [19–22]. For example,  $\text{La}_2\text{Zr}_3(\text{MoO}_4)_9$  ceramic sintered at 775 °C was reported with good microwave dielectric properties of  $\epsilon_r = 10.8$ ,  $Qf = 50,628$  GHz and  $\tau_f = -38.8$  ppm/°C [20]. However, a new  $\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$  ceramic have not been reported up to now. Therefore, the CZM samples were

prepared successfully via the solid-state route for the first time. The microstructures, sintering characteristics and microwave dielectric properties of CZM were studied scientifically. The chemical bonds theory of complex crystals and infrared (IR) reflectivity spectra were helpful to explore the relationships between the intrinsic factors and dielectric properties of CZM samples.

## 2. Experimental procedure

The CZM samples were prepared by the standard solid-state reaction route. High purity raw materials of cerium oxide ( $\text{CeO}_2$ , 99.9%, Macklin), zirconium dioxide ( $\text{ZrO}_2$ , 99.99%, Macklin), and molybdenum trioxide ( $\text{MoO}_3$ , 99.95%, Nine-Dinn chemistry Shanghai Co., Ltd, China) were accurately weighed in  $\text{CeO}_2$ – $\text{ZrO}_2$ – $\text{MoO}_3$  system with molar ratio 2:3:9. The mixed powders were placed in nylon pot with the spherical  $\text{ZrO}_2$  ceramics and anhydrous alcohol as grinding medium, and then the ball-machine was set to rotate at 200 r/min for ball-milling for 24 h. The obtained slurries separated from zirconia spheres were dried in an open oven and then calcined at 700 °C for 2 h. The pre-sintered powders were re-grounded for 24 h, and dried at 80 °C once more. Afterwards, the granulated that added with 10 % PVA as an adhesive were pressed into cylinders sample under a uniaxial pressure of 200 MPa. At last, the prepared CZM disks were sintered at

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**Table 1**  
Summarized microwave dielectric properties of Mo-based microwave dielectric ceramics.

Materials	Sintering temperatures (°C)	$\epsilon_r$	$Qf$ (GHz)	$\tau_f$ (ppm/°C)	References
$\text{Sm}_2\text{Zr}_3(\text{MoO}_4)_9$	875	11.0	74,012	−45.3	[19]
$\text{Nd}_2\text{Zr}_3(\text{MoO}_4)_9$	850	10.8	58,942	−40.9	[19]
$\text{La}_2\text{Zr}_3(\text{MoO}_4)_9$	775	10.8	50,628	−38.8	[20]
$\text{La}_2\text{Zr}_3(\text{MoO}_4)_9$	650	10.8	61,790	−29.1	[21]
$\text{La}_2(\text{Zr}_{0.92}\text{Ti}_{0.08})_3(\text{MoO}_4)_9$	750	10.33	80,658	3.48	[21]
$\text{Eu}_2\text{Zr}_3(\text{MoO}_4)_9$	600	10.75	74,900	−8.88	[22]
$\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$	575	10.69	19,062	−1.29	This work

500–625 °C for 4 h under air atmosphere.

The phase structures of the sintered CZM ceramics were characterized via an X-ray diffractometer (XRD: D8 Advance, Bruker Co., Germany) with nickel-filtered Cu-K $\alpha$  radiation ( $\lambda = 1.542 \text{ \AA}$ , 40 kV/40 mA). The FullProf\_Suite program [23] was employed to obtain the unit cell parameters from the Rietveld refinement of XRD data. The sintered specimens were magnified appropriate times to observe the surface micro-structures by using QUANTA 250FEG type Scanning electron microscope (SEM: FEI Co., United States) coupled with EDS. The room-temperature infrared reflectivity spectrum data of thin polished disc (around at 1 mm) were taken using a Fourier transform infrared spectrometer (FTIR: Bruker IFS 66v) at National Synchrotron Radiation Laboratory (NRS), BL01B infrared beamline station, University of Science and Technology of China). The dielectric constants were measured at approximately 12.5–13.5 GHz in the light of Hakki-Coleman dielectric resonator method [24]. The unload quality factors were measured with the help of TE<sub>018</sub> cavity method [25] using TE01d mode where a pellet was placed in metal shielded cavity to determine the coupling energy of the entire system through the coaxial probe. The temperature coefficient of resonant frequency values were attained by using Eq. (1):

$$\tau_f = \frac{f_T - f}{60 \times f_0} \times 10^6 (\text{ppm}/^\circ\text{C}) \quad (1)$$

where,  $f_T$  and  $f_0$  were the resonant frequency at 85 °C and 25 °C, respectively.

The apparent densities of as-fired samples were gauged through liquid Archimedeian method (FR124N electronic balance, Ohaus Instruments Shanghai Co., Ltd.) utilizing deionized water as liquid media. The theoretical densities ( $\rho_{\text{theory}}$ ) were estimated by formula (2):

$$\rho_{\text{theory}} = \frac{ZA}{V_c N_A} (\text{g}/\text{cm}^3) \quad (2)$$

where  $V_c$  was volume of refined unit-cell ( $\text{cm}^3$ ),  $N_A$  was Avogadro constant ( $6.023 \times 10^{23} \text{ mol}^{-1}$ ),  $Z$  was atomic unit cell number ( $Z = 6$  in this study),  $A$  represented the relative atomic weight ( $\text{g}/\text{mol}^{-1}$ ). The following equation was applied to reckon the relative density ( $\rho_{\text{relative}}$ ):

$$\rho_{\text{relative}} = \frac{\rho_{\text{apparent}}}{\rho_{\text{theory}}} \times 100 \% \quad (3)$$

### 3. Results and discussion

Fig. 1 simultaneously presented the typical XRD patterns of CZM sample sintered at 500–625 °C for 4 h and the calcined powders sintered at 700 °C for 2 h. It could be seen that the XRD diffraction peak positions of all experimental specimens matched well with  $\text{Pr}_2\text{Zr}_3(\text{MoO}_4)_9$ -like crystal structure. And there were no traces of the additional impurity phases during the sintering process, which indicated that the pure single trigonal solid-solution phase structure with a R-3c (167) space group was formed in all samples.

In view of Rietveld refinement method (Fig. 2), the refined lattice parameters of the samples sintered within the temperature range of 500–625 °C from XRD data were gained so as to better investigate the

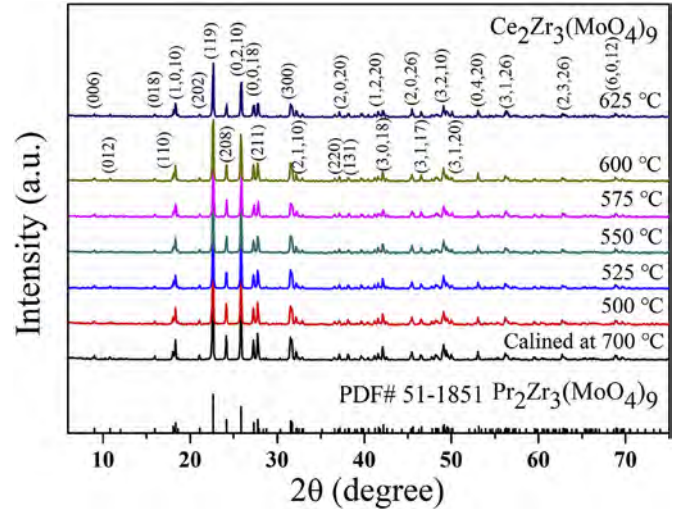


Fig. 1. XRD patterns of  $\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$  ceramics sintered at 500–625 °C.

CZM structure at great length, where  $\text{Nd}_2\text{Zr}_3(\text{MoO}_4)_9$  (ICSD File No. 92,600) served as initial structural model. Fig. 3 depicted the variation tendencies of lattice parameters. The  $a$  and  $b$  gradually increased with the increasing sintering temperature, and  $c$  was roughly stable at  $58.8629 \text{ \AA}$ , which resulted in the increase of  $V_m$  from  $4929.10 \text{ \AA}^3$  at 500 °C to  $4932.40 \text{ \AA}^3$  at 625 °C. The structural parameters and refined discrepancy factors ( $R_{wp}$ ,  $R_p$ , and  $\chi^2$ ) after Rietveld refinement were all listed in Table 2.  $R_p$ ,  $R_{wp}$  and  $\chi^2$  were measured as 7.08–8.73%, 9.20–11.0% and 1.57–2.09, respectively, which confirmed the phase purity of CZM.

The schematic illustration of trigonal structured CZM sample was shown in Fig. 4, and Table 3 exhibited the occupations and the refined atomic positions of CZM sample. The CZM crystal was constituted of  $\text{CeO}_9$ ,  $\text{ZrO}_6$  and  $\text{MoO}_4$  polyhedra, which were shared by vertices. In this structure, Ce, Zr(1), Zr(2), Mo(1), Mo(2) and O atoms occupied the 12c, 6b, 12c, 36f, 18e and 36f Wyckoff positions, respectively.

Fig. 5 illustrated the apparent and relative densities of CZM samples sintered at 500–625 °C. When the temperature increased to 575 °C, apparent density increased from 3.59 to the maximum value of 3.84  $\text{g}/\text{cm}^3$ , and then reached a saturated value of 3.83  $\text{g}/\text{cm}^3$  at 575–625 °C. The increase in density was due to the decrease of porosity. Just like the trend of apparent density, the relative density increased from 88.99 % at 500 °C to 95.30 % at 575 °C and then kept stable. SEM micrographs sintered at 500–625 °C for 4 h were shown in Fig. 6. Slight pores were observed at a lower temperature of 500 °C, and dense microstructures with clear grain boundaries could be obtained as shown in Fig. 6(b)–(f). EDS of CZM sample sintered at 575 °C was shown in Fig. 6(g). Atom ratio of Ce, Zr, Mo and O were 2.85 %, 5.43 %, 15.00 % and 76.73 % respectively, which was consistent with chemical formula of CZM.

Microwave dielectric properties for CZM samples sintered at 500–625 °C were presented in Fig. 7. With temperature increasing, dielectric constant increased from 9.51 at 500 °C to 10.69 at 575 °C, and then saturated at 575–625 °C. It was obvious that the variation in

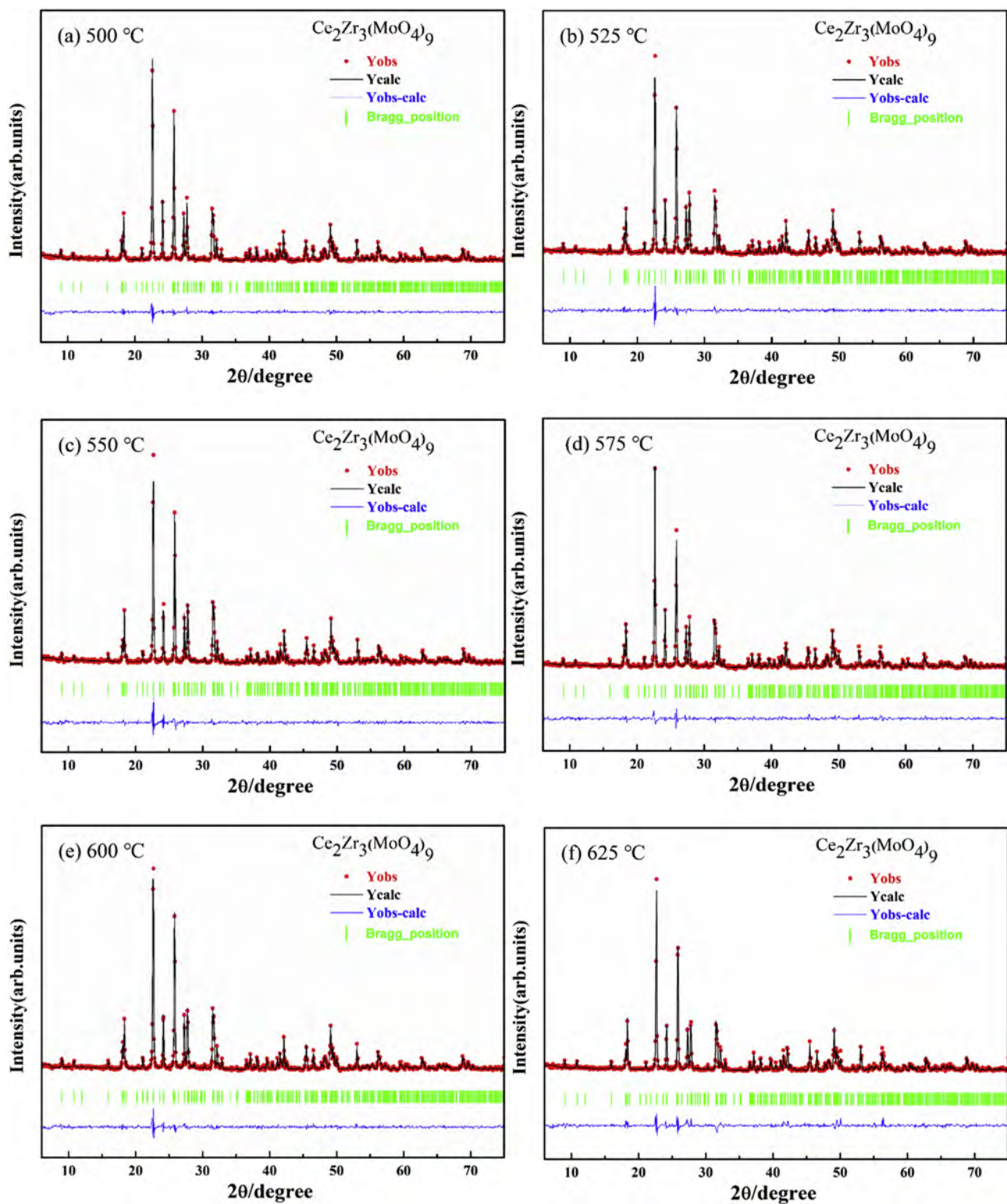


Fig. 2. Rietveld refinement of  $\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$  ceramics sintered at different temperatures (a-f corresponding to 500–625 °C).

dielectric constant was consistent with the density. The higher density corresponded to higher dielectric constant. The dielectric constant of ceramic with single phase is mostly related to the density and polarizability. To further investigate the effect of polarizability on dielectric constant, molecular polarizability of CZM sample is calculated according to Shannon's report based on Eq. (4) [26].

$$\alpha_{\text{theo.}} = 2\alpha(\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9) = 2\alpha((\text{Ce}^{3+}) + 3\alpha(\text{Zr}^{4+}) + 9\alpha(\text{Mo}^{6+}) + 36\alpha(\text{O}^{2-})) \quad (4)$$

where  $\alpha$  of  $\text{Ce}^{3+}$ ,  $\text{Zr}^{4+}$  and  $\text{O}^{2-}$  are the theoretical ion polarizabilities reported by Shannon et al. [26] and  $\alpha(\text{Mo}^{6+})$  is calculated by Choi et al. [27]. The observed value dielectric polarizability is obtained by using Clausius-Mosotti equation (5):



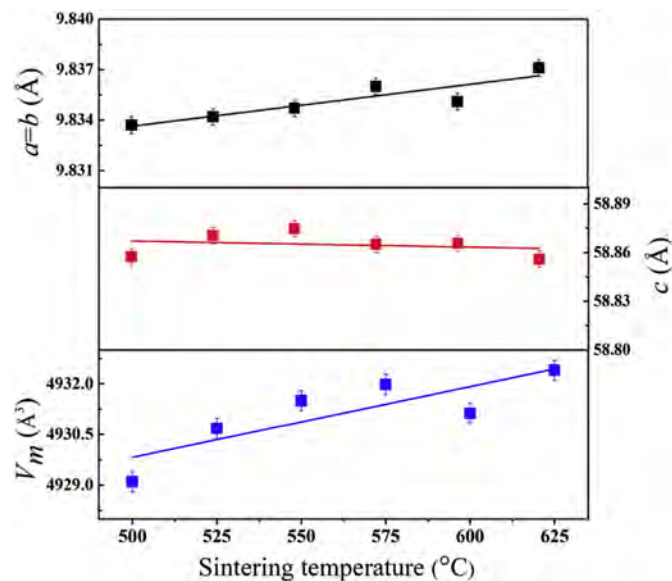


Fig. 3. The variation tendencies of lattice parameters with sintering temperature.

Table 2

The refinement parameters, theoretical densities and relative densities of  $\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$  ceramics fired at 500–625 °C.

Structural parameters	sintering temperature (°C)					
	500	525	550	575	600	625
$a = b$ (Å)	9.8337	9.8342	9.8347	9.8360	9.8351	9.8371
$c$ (Å)	58.8574	58.8705	58.8747	58.8650	58.8656	58.8559
$\alpha = \beta$ (°)	90	90	90	90	90	90
$\gamma$ (°)	120	120	120	120	120	120
$V_m$ (Å <sup>3</sup> )	4929.10	4930.68	4931.50	4931.98	4931.13	4932.40
$R_p$ (%)	7.08	7.26	7.53	7.53	7.30	8.73
$R_{wp}$ (%)	9.20	9.33	9.76	9.86	9.38	11.0
$\chi^2$	1.61	1.57	1.68	1.67	1.66	2.09
$\rho_{theo.}$ (g·cm <sup>-3</sup> )	4.0305	4.0293	4.0286	4.0282	4.0289	4.0279
$\rho_{rela.}$ (%)	88.99	92.28	94.14	95.30	95.07	94.96

$\rho_{thep.}$ — theoretical density of  $\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$  ceramics.

$P_{rela.}$ — relative density of  $\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$  ceramics.

$R_p$ — reliability factor of patterns.

$R_{wp}$ — the reliability factor of weighted patterns.

$\chi^2$ — goodness of fit indicator =  $\chi^2 = (R_{wp}/R_{exp})^2$ .

Table 3

Atomic fractional coordinates of the  $\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$  ceramics fired at 575 °C.

Atom	Wyckoff position	Site	x	y	z	Occupancy
Ce	12c	3	0.6667	0.3333	0.0297	0.3333
Zr1	6b	−3	0.0000	0.0000	0.0000	0.1667
Zr2	12c	3	0.3333	0.6667	0.0707	0.3333
Mo1	36f	1	0.0979	0.7284	0.0248	1.0000
Mo2	18e	2	0.2874	0.2874	0.2500	0.5000
O1	36f	1	0.1056	0.5876	0.0311	1.0000
O2	36f	1	0.1577	0.6457	0.0035	1.0000
O3	36f	1	0.2008	0.7397	0.0502	1.0000
O4	36f	1	0.1277	0.9267	0.0182	1.0000
O5	36f	1	0.2887	0.4747	0.2467	1.0000
O6	36f	1	0.1647	0.1926	0.2731	1.0000

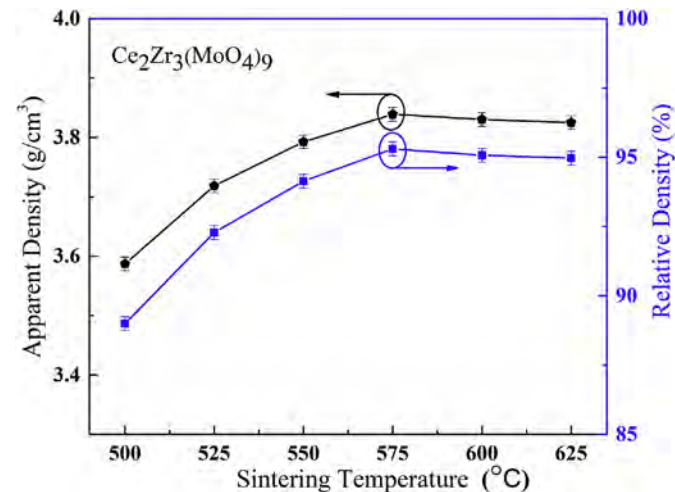


Fig. 5. Apparent and relative density of  $\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$  ceramics as a function of sintering temperature.

$$\alpha_{obs} = \frac{1}{b} V_m \frac{\epsilon_r - 1}{\epsilon_r + 2} \quad (5)$$

where  $V_m$  is the molar volume and  $b$  is a constant value ( $4\pi/3$ ). The theoretical dielectric polarizability was about 123.93, which was closed to the observed value (149.85), suggesting the accuracy of the theoretical calculation of the dielectric polarizability.

In addition, the quality factor increased from 15,230 GHz at 500 °C to 19,975 GHz at 525 °C, and then kept stable at 525–575 °C. The

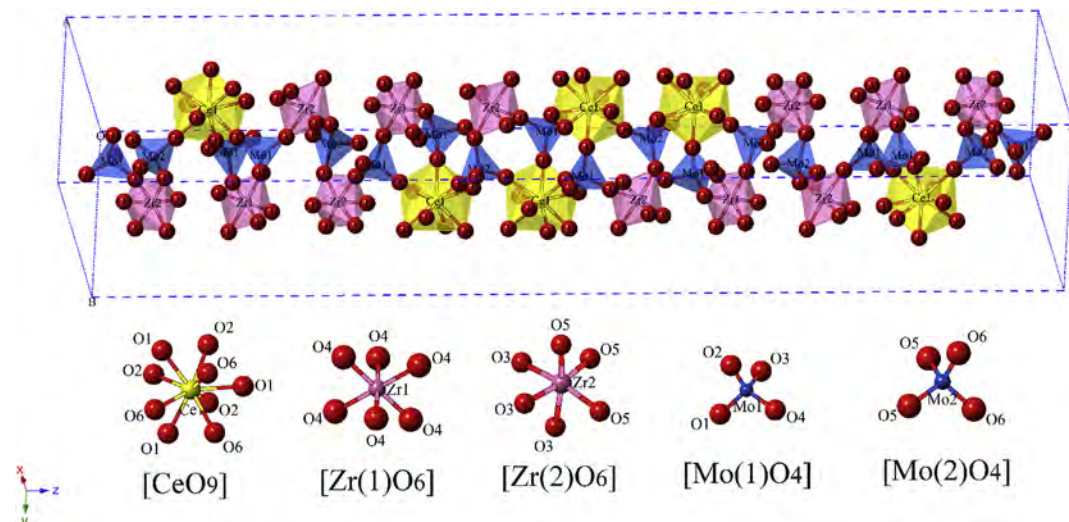


Fig. 4. Schematic diagram of the crystal structure for  $\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$  ceramics.

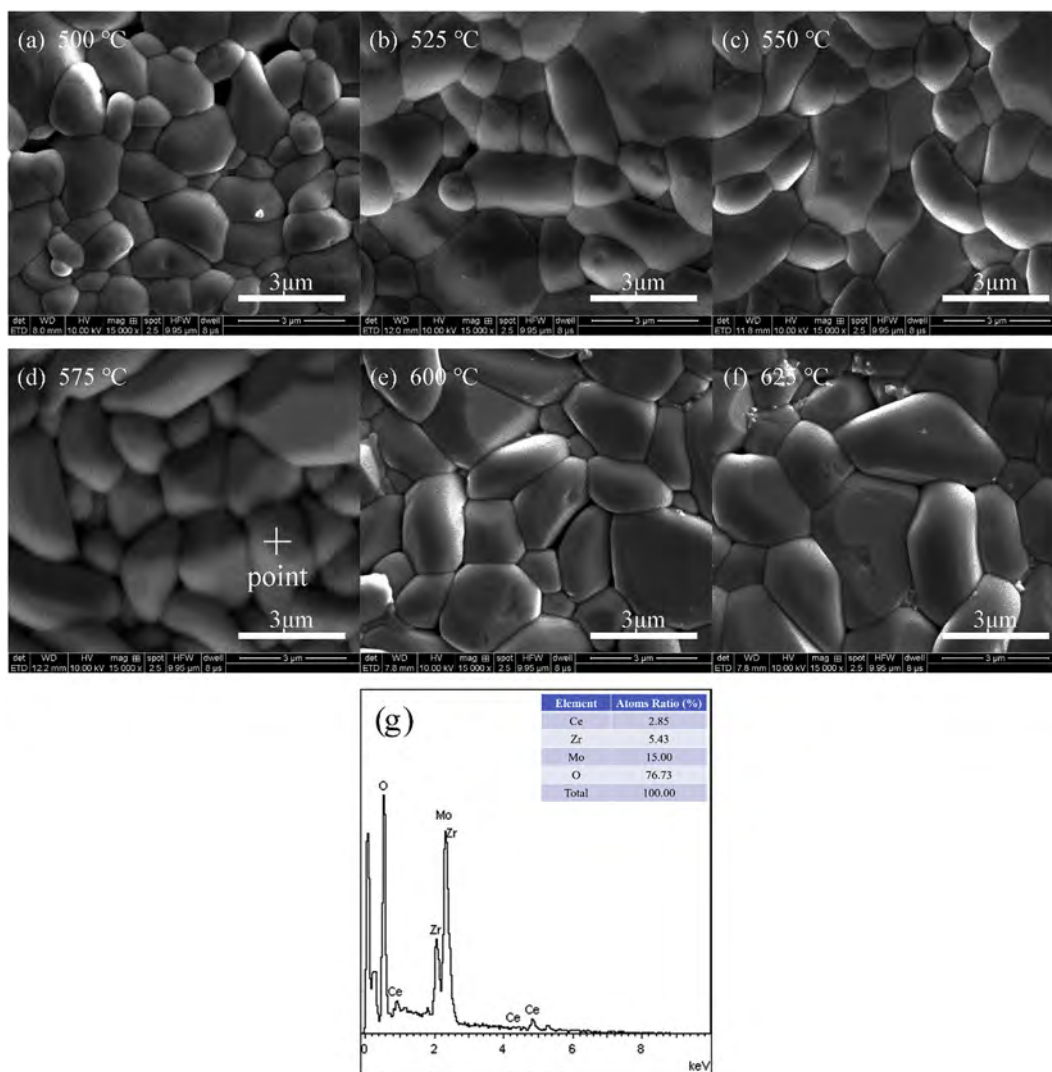


Fig. 6. Typical SEM micrographs of  $\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$  ceramics fired at (a) 500 °C, (b) 525 °C, (c) 550 °C, (d) 575 °C, (e) 600 °C, (f) 625 °C and (g) EDS analysis about grains selected randomly from the sample sintered at 575 °C.

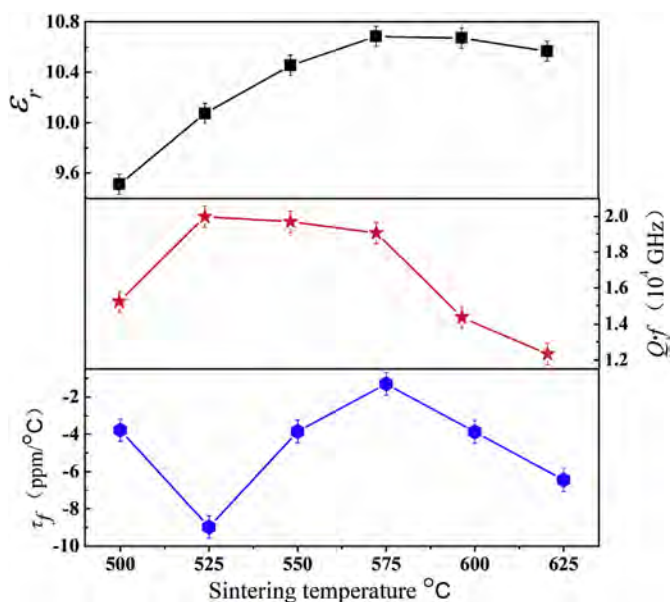


Fig. 7. Curves of  $\epsilon_r$ ,  $Qf$ , and  $\tau_f$  values of the  $\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$  ceramics sintered at 500–625 °C.

quality factor declined slightly with further increase of sintering temperature. With an increase in temperature from 500 to 625 °C, the saturation  $\tau_f$  value of  $-4.7$  ppm/°C was obtained. In general, the microwave dielectric properties are associate with the extrinsic factors, including density, the secondary phase, grain size, impurities, pores, etc., and intrinsic factors, dependent on the structure characteristics [28]. The influences of extrinsic factors on the properties can be neglected when the sintered samples are compact. The relationships between dielectric properties and crystal structure should be investigated. Chemical bonds theory of complex crystals involving bond ionicity  $f_b$ , lattice energy  $U$ , bond energy  $E$ , coefficient of thermal expansion  $\alpha$  had been studied to explore the effects of chemical bond on crystal structure. Therefore, the influences of chemical bonds on microwave dielectric properties could be obtained. The sum of binary crystals could be written as shown in Eq. (6) based on the Phillips-Van Vechten-Levine (P-V-L) theory [29–31]. Effective valence of Ce, Zr and Mo were  $P_{\text{Ce}} = 3$ ,  $P_{\text{Zr}} = 4$  and  $P_{\text{Mo}} = 6$ , respectively. The effective valence of O should be followed by the subformula.

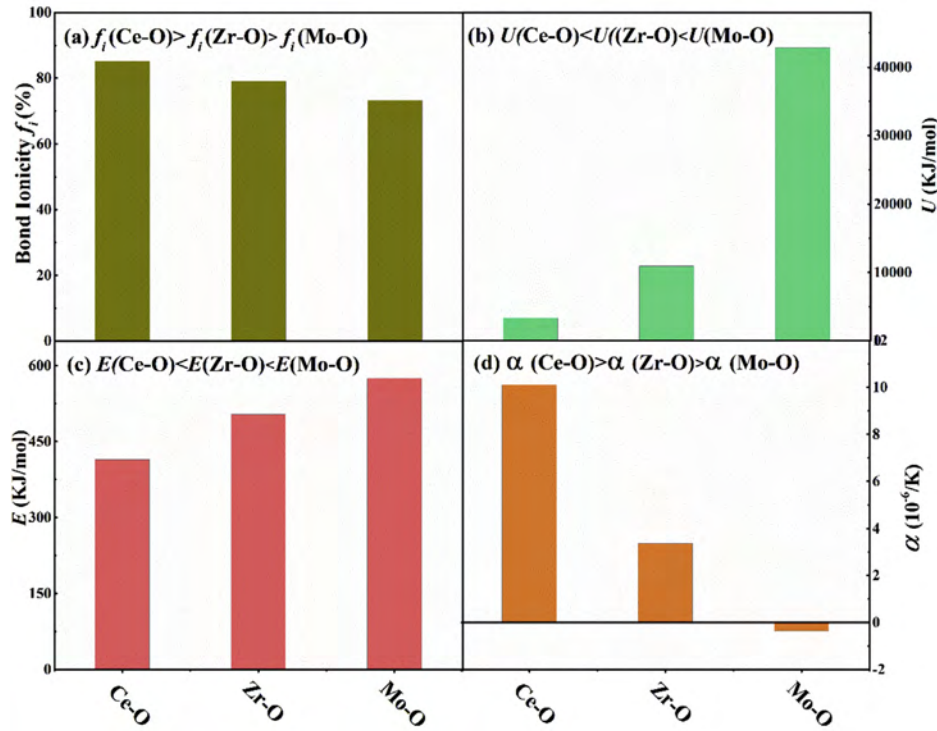
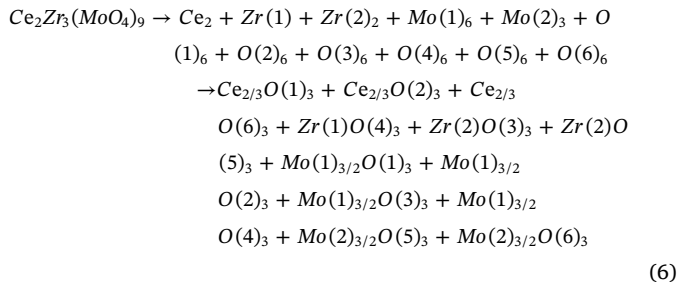


Fig. 8. (a) Bond ionicity, (b) lattice energy, (c) bond energy, and (d) thermal expansion coefficient for  $Ce_2Zr_3(MoO_4)_9$  ceramics sintered at 575 °C.



The dielectric constant can be predicted based on the  $f_i$  value as shown in Fig. 8(a) [32]. The bond ionicity can be calculated by using Eqs. (7–11) [29–31]:

$$\epsilon = \frac{n^2 - 1}{1 - f_i} + 1 \quad (7)$$

$$f_i^\mu = \frac{(C^\mu)^2}{(E_g^\mu)^2} \quad (8)$$

$$(E_g^\mu)^2 = (E_h^\mu)^2 + (C^\mu)^2 \quad (9)$$

$$(E_h^\mu)^2 = \frac{39.74}{(d^\mu)^{2.48}} \quad (10)$$

$$C^\mu = 14.4b^\mu \exp(-k_s^\mu r_0^\mu) \left[ (Z_A^\mu)^* - \frac{n}{m} (Z_B^\mu)^* \right] / r_0^\mu \quad (11)$$

where  $E_g^\mu$  is the average energy gap,  $E_h^\mu$  is the homopolar part,  $C^\mu$  is the heteropolar part and  $\exp(-k_s^\mu r_0^\mu)$  is thomas-Fermi screening factor [33]. The average ionicity of Ce–O, Zr–O and Mo–O bonds were 85.18 %, 79.08 % and 73.24 %, respectively. Therefore, the  $f_i$  value of Ce–O bond provided the majority contribution to the  $\epsilon_r$  value of CZM sample.

The lattice energy of CZM sample is evaluated by Eq. (12–15) [29–31], which can be decomposed into ionic part  $U_{bi}^\mu$  and covalent part  $U_{bc}^\mu$ :

$$U_{cal} = \sum_{\mu} U_b^\mu \quad (12)$$

$$U_b^\mu = U_{bc}^\mu + U_{bi}^\mu \quad (13)$$

$$U_{bc}^\mu = 2100m \frac{(Z_A^\mu)^{1.64}}{(d^\mu)^{0.75}} f_c^\mu \quad (14)$$

$$U_{bi}^\mu = 1270 \frac{(m+n)Z_A^\mu Z_B^\mu}{d^\mu} \left( 1 - \frac{0.4}{d^\mu} \right) f_i^\mu \quad (15)$$

The maximum average lattice energies was 42,856 kJ/mol, which was obtained in Mo–O bond. It could be concluded that Mo–O bond played a predominant role in affecting quality factor of CZM sample.

Moreover, the bond energy measures the strength of chemical bonds, which is calculated based on the following formulas [34–36]:

$$E^\mu = t_c E_c^\mu + t_i E_i^\mu \quad (16)$$

$$E_c^\mu = \frac{(r_{cA} + r_{cB})}{d^\mu} (E_{A-A} E_{B-B})^{1/2} (KJ/mol) \quad (17)$$

$$E_i^\mu = \frac{1389.088}{d^\mu} (KJ/mol) \quad (18)$$

$$t_i = \left| \frac{(S_A - S_B)}{6} \right| \quad (19)$$

$$t_c + t_i = 1 \quad (20)$$

where  $S_A$  and  $S_B$  are the electronegativities of A and B ions;  $r_{cA}$  and  $r_{cB}$  are the covalent radii;  $t_i$  and  $t_c$  are ionic and covalent proportional coefficient of an individual bond  $\mu$ ; homonuclear bond energy  $E_{A-A}$  and  $E_{B-B}$  can be consulted in handbook of bond energies [37]. According to Fig. 8(c) and Table 4, the sequence of  $E(Mo-O) > E(Zr-O) > E(Ce-O)$  suggested that Mo–O bond provided a major contribution to CZM structure.

The correlation among the temperature coefficient of resonant frequency  $\tau_f$ , temperature coefficient of the dielectric constant  $\tau_\epsilon$  and coefficient of thermal expansion  $\alpha$  is shown in Eq. (21), in which  $\alpha$  is

**Table 4**Bond ionicity, bond energy, lattice energy, and thermal expansion coefficient of  $\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$  ceramic sintered at 575 °C.

Bond type	$d$ (Å)	$f_l^u$ (%)	$f_c^u$ (%)	$U_{bc}$ (kJ/mol)	$U_{bi}$ (kJ/mol)	$U$ (kJ/mol)	$E_i$ (kJ/mol)	$E_c$ (kJ/mol)	$E$ (kJ/mol)	$\alpha$ ( $10^{-6}\text{K}^{-1}$ )
Ce–O(1) <sup>1</sup>	2.381(2)	84.61	15.39	681	2753	3434	336.144(3)	583.077(4)	431.625(1)	9.641(7)
Ce–O(1) <sup>2</sup>	2.381(8)	84.61	15.39	681	2753	3434	336.059(6)	582.930(6)	431.516(4)	9.641(7)
Ce–O(1) <sup>3</sup>	2.382(2)	84.61	15.39	681	2753	3434	336.003(2)	582.832(7)	431.443(9)	9.641(7)
Ce–O(2) <sup>1</sup>	2.549(9)	85.54	14.46	608	2634	3242	313.905(2)	544.501(4)	403.069(0)	10.400(3)
Ce–O(2) <sup>2</sup>	2.550(1)	85.54	14.46	608	2634	3242	313.880(6)	544.458(6)	403.037(4)	10.400(3)
Ce–O(2) <sup>3</sup>	2.550(5)	85.54	14.46	608	2634	3242	313.831(3)	544.373(3)	402.974(2)	10.400(3)
Ce–O(6) <sup>1</sup>	2.520(1)	85.39	14.61	620	2655	3275	317.617(1)	550.940(0)	407.835(3)	10.263(6)
Ce–O(6) <sup>2</sup>	2.520(4)	85.39	14.61	620	2655	3275	317.579(3)	550.874(5)	407.786(8)	10.263(6)
Ce–O(6) <sup>3</sup>	2.520(8)	85.39	14.61	620	2654	3274	317.528(9)	550.787(1)	407.722(0)	10.267(7)
Zr(1)–O(4) × 6	2.036(5)	79.04	20.96	2508	8450	10,958	410.771(9)	681.769(7)	506.072(8)	3.344(4)
Zr(2)–O(3) <sup>1</sup>	2.147(9)	79.92	20.08	2309	8204	10,513	389.467(4)	646.410(0)	479.825(5)	3.620(1)
Zr(2)–O(3) <sup>2</sup>	2.148(5)	79.92	20.08	2308	8202	10,510	389.358(6)	646.229(5)	479.691(5)	3.622(0)
Zr(2)–O(3) <sup>3</sup>	2.148(7)	79.93	20.07	2307	8202	10,509	389.322(4)	646.169(3)	479.646(9)	3.622(7)
Zr(2)–O(5) <sup>1</sup>	1.955(3)	78.32	21.68	2674	8632	11,306	427.830(5)	710.082(3)	527.089(1)	3.144(0)
Zr(2)–O(5) <sup>2</sup>	1.955(6)	78.33	21.67	2674	8632	11,306	427.764(9)	709.973(4)	527.008(2)	3.144(0)
Zr(2)–O(5) <sup>3</sup>	1.956(2)	78.33	21.67	2672	8630	11,302	427.633(7)	709.755(6)	526.846(6)	3.146(2)
Mo(1)–O(1)	1.810(6)	73.38	26.62	10,148	32,479	42,627	517.177(0)	766.830(9)	570.436(5)	–0.343(2)
Mo(1)–O(2)	1.750(1)	72.77	27.23	10,645	33,000	43,645	535.055(5)	793.339(8)	590.156(2)	–0.409(1)
Mo(1)–O(3)	1.773(9)	73.02	26.98	10,444	32,795	43,239	527.876(8)	782.695(8)	582.238(2)	–0.383(2)
Mo(1)–O(4)	1.858(8)	73.82	26.18	9784	32,062	41,846	503.766(2)	746.9464	555.644(7)	–0.290(5)
Mo(2)–O(5) × 2	1.842(7)	73.68	26.32	9903	32,202	42,105	508.167(7)	753.472(6)	560.499(4)	–0.308(2)
Mo(2)–O(6) × 2	1.750(4)	72.78	27.22	10,643	32,997	43,640	534.963(8)	793.203(8)	590.055(0)	–0.408(8)
Ce–O <sub>avg</sub>	2.484(1)	85.18	14.82	-	-	3317	-	-	414.112(2)	10.102(3)
Zr–O <sub>avg</sub>	2.044(3)	79.08	20.92	-	-	10,933	-	-	503.740(0)	3.363(8)
Mo–O <sub>avg</sub>	1.797(5)	73.24	26.76	-	-	42,856	-	-	574.838(3)	–0.357(5)

calculated from Eq. (22–25):

$$\tau_f = -\left(\frac{\tau_c}{2} + \alpha\right) \quad (21)$$

$$\alpha = \sum_{\mu} F_{mn}^{\mu} \alpha_{mn}^{\mu} \quad (22)$$

$$\alpha_{mn}^{\mu} = -3.1685 + 0.8376\gamma_{mn} \quad (23)$$

$$\gamma_{mn} = \frac{kZ_A^{\mu}N_{CA}^{\mu}}{U_b^{\mu}\Delta_A}\beta_{mn} \quad (24)$$

$$\beta_{mn} = \frac{m(m+n)}{2n} \quad (25)$$

where  $F_{mn}^{\mu}$  is the proportion of  $\mu$  bond,  $N_{CA}^{\mu}$  is the coordination number of  $\mu$  bond for cation,  $Z_A^{\mu}$  is the valence states of cation, and  $k$  is Boltzmann constant. The average  $\alpha$  of Ce–O, Zr–O and Mo–O were  $10.10 \times 10^{-6}/\text{K}$ ,  $3.36 \times 10^{-6}/\text{K}$  and  $-0.36 \times 10^{-6}/\text{K}$ , respectively. Hence, the main positive effect to temperature coefficients of resonant frequency was the  $\alpha$  value of Mo–O bond because of  $\alpha_{\text{Mo-O}} < 0$ .

Fig. 9(a) presented the IR reflectivity spectrum of CZM sintered at 575 °C. Eq. (26) is used to calculate the complex dielectric permittivity  $\varepsilon^*(\omega)$  based on the model of classical harmonic oscillator, and the complex reflectivity  $R$  can be obtained as Eq. (27) [38–40]:

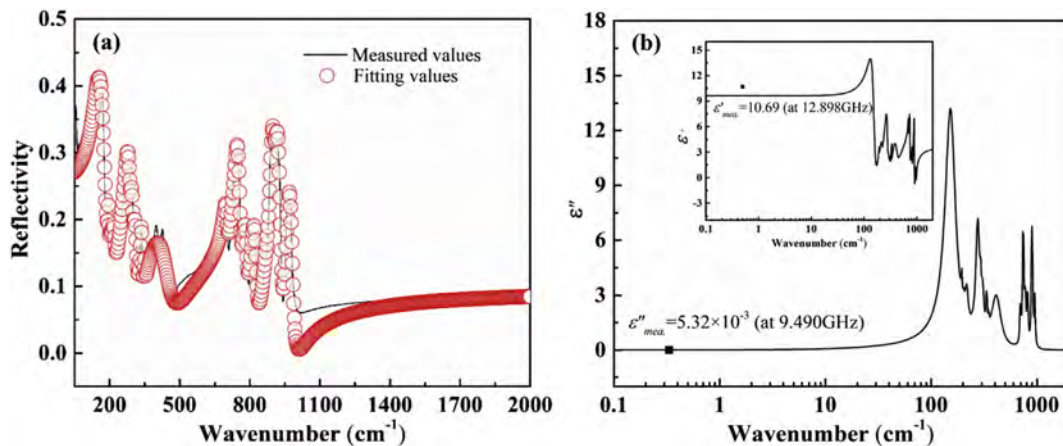
$$\varepsilon^*(\omega) = \varepsilon_{\infty} + \sum_{j=1}^n \frac{\omega_{pj}^2}{\omega_{oj}^2 - \omega^2 + j\omega\gamma_j} \quad (26)$$

$$R = \left| \frac{1 - \sqrt{\varepsilon^*(\omega)}}{1 + \sqrt{\varepsilon^*(\omega)}} \right|^2 \quad (27)$$

where  $\omega_{oj}$ ,  $\omega_{pj}$  and  $\varepsilon_{\infty}$  are the the eigen frequency, plasma frequency and dielectric constant, respectively;  $\gamma_j$  is damping factor;  $n$  is the number of transverse phonon modes. In addition, the following formula were used to calculate the dielectric loss tangent  $\tan \delta$ :

$$\tan \delta = \frac{\varepsilon''}{\varepsilon'} = \frac{\sum_{j=1}^n \Delta\varepsilon_j(\gamma_j\omega)/\omega_{oj}^2}{\varepsilon_{\infty} + \sum_{j=1}^n \Delta\varepsilon_j} \quad (28)$$

The complex permittivity values and fitted IR reflectivity spectra



**Fig. 9.** (a) Measured (black line) and fitted (red line) infrared reflectivity spectrum, and (b) real and imaginary parts of the complex permittivity for  $\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$  ceramics sintered at 575 °C.



**Table 5**

Phonon parameters obtained from the fitting of the infrared spectra of  $\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$  ceramic.

Mode	$\text{Ce}_2\text{Zr}_3(\text{MoO}_4)_9$ $\epsilon_\infty = 3.55$			
	$\omega_{oj}$	$\omega_{pj}$	$\gamma_j$	$\Delta\epsilon_j$
1	154.01	291.42	42.58	3.58
2	194.74	34.04	4.491	0.03
3	203.99	29.79	7.42	0.02
4	217.64	84.17	17.87	0.15
5	276.80	217.31	27.34	0.62
6	294.71	66.58	9.66	0.05
7	308.27	103.64	18.15	0.11
8	335.93	75.77	7.54	0.05
9	351.08	36.62	4.84	0.01
10	417.08	348.80	103.06	0.70
11	700.34	170.43	16.36	0.06
12	744.83	360.40	24.58	0.23
13	779.66	205.33	21.30	0.07
14	817.31	254.14	25.03	0.10
15	895.36	327.63	17.93	0.13
16	911.61	287.55	26.02	0.10
17	960.27	251.04	24.66	0.07

$\omega_{oj}$ ,  $\omega_{pj}$ ,  $\gamma_j$  and  $\Delta\epsilon_j$  are the eigen frequency, plasma frequency, coefficient of damping and calculated  $\epsilon$  of individual mode,  $\epsilon_\infty$  is the dielectric constant affected by electronic polarization in higher frequencies.

were exhibited in Fig. 9(b). There were 17 internal modes based on Table 5. The extra-polated permittivity of CZM was 9.63, which was close to the measured value (10.69). In addition, the calculated dielectric loss of  $2.27 \times 10^{-3}$  was lower than the experimental value of  $5.32 \times 10^{-3}$ , which attributed to various extrinsic loss. Hence, the most polarization contribution for CZM was mainly related to the absorptions of phonon oscillation at infrared region.

#### 4. Conclusion

A new CZM ceramic was fabricated using the solid-state reaction method at a low sintering temperature. A single phase, with trigonal structure in R-3c space group, was confirmed by the X-ray diffractometer. Rietveld refinements were used to further investigate the crystal structure. The dense samples sintered at 575–625 °C with homogeneous microstructures were characterized by Scanning electron microscopy. The structure-property relations of CZM samples were studied by using the chemical bonds theory of complex crystals and far-infrared spectrum. The optimum microwave dielectric properties with  $\epsilon_r = 10.69$ ,  $Qf = 19,062$  GHz and  $\tau_f = -1.29$  ppm/°C were achieved in CZM sample sintered at 575 °C.

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